

# Hydrogen crossover in high-temperature PEM fuel cells

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## Abstract

In this paper, hydrogen crossover was measured in an environment of high-temperature proton exchange membrane (PEM) fuel cells using a steady-state electrochemical method at various temperatures ( $T$ ) (80–120 °C), backpressures ( $P$ ) (1.0–3.0 atm), and relative humidities (RH) (25–100%). An H<sub>2</sub> crossover model based on an MEA consisting of five layers – anode gas diffusion layer, anode catalyst layer, proton exchange membrane (Nafion 112 or Nafion 117), cathode catalyst layer, and cathode gas diffusion layer – was constructed to obtain an expression for H<sub>2</sub> permeability coefficients as a function of measured H<sub>2</sub> crossover rates and controlled H<sub>2</sub> partial pressures. The model analysis suggests that the dominant factor in the overall H<sub>2</sub> crossover is the step of H<sub>2</sub> diffusing through the PEM. The H<sub>2</sub> permeability coefficients as a function of  $T$ ,  $P$ , and RH obtained in this study show that the increases in both  $T$  and  $P$  could increase the H<sub>2</sub> permeability coefficient at any given RH. However, the effect of RH on the permeability coefficient seems to be more complicated. The  $T$  effect is much larger than that of  $P$  and RH. Through experimental data simulation an equation was obtained to describe the  $T$  dependencies of the H<sub>2</sub> permeability coefficient, based on which other parameters such as maximum permeability coefficients and activation energies for H<sub>2</sub> crossover through both Nafion 112 and 117 membranes were also evaluated. Both Nafion 112 and Nafion 117 showed similar values of such parameters, suggesting that membrane thickness does not play a significant role in the H<sub>2</sub> crossover mechanism.

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## 1. Introduction

Proton exchange membrane (PEM) fuel cells, used as a power source in transportation and stationary applications, have been recognized as the most promising energy converters in terms of high power density, low or zero emissions, and high efficiency [1]. In the last several decades, in order to realize the commercialization of PEM fuel cells, tremendous effort has been put into improving performance and addressing several technology gaps, including high cost and limited reliability and durability. With respect to this effort, property improvements in the membrane electrode assembly (MEA), the heart of a fuel cell in which the converting reactions occur, are the major focus.

An MEA is made up of a cathode, an anode, and a PEM. Their impedances are primarily responsible for the performance loss that occurs during fuel cell operation. For example, fuel cell polarization can be expressed as Eq. (1) [2]

$$V_{\text{cell}} = E^{\text{OCV}} - \eta_c - \eta_a - I_{\text{cell}} R_m \quad (1)$$

where  $E^{\text{OCV}}$  is the fuel cell open circuit voltage (OCV),  $I_{\text{cell}}$  the fuel cell current density,  $R_m$  the MEA resistance dominated by the membrane resistance, and  $\eta_c$  and  $\eta_a$  are the overpotentials for the cathode and anode, respectively. Among the voltage losses of the cathode ( $\eta_c$ ), anode ( $\eta_a$ ), and membrane ( $I_{\text{cell}} R_m$ ), that of the membrane ( $I_{\text{cell}} R_m / V_{\text{cell}}$ ) accounts for about 30% [3]. Reduction in membrane proton resistance is the major approach to reducing membrane voltage loss. One effective and straightforward way to reduce membrane resistance is by using a thinner membrane in the MEA. For example, in early fuel cell development, Nafion 117 with a thickness of 175–183  $\mu\text{m}$  was widely used. Later, Nafion 115 with a thickness of 127  $\mu\text{m}$  was employed, followed by Nafion 112 with a thickness of 51  $\mu\text{m}$  [4], and most recently,

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Nafion 211 or Nafion 111 with a thickness of 25  $\mu\text{m}$ . However, with decreasing the membrane thickness, another problem, fuel (typically, hydrogen) crossover, may become a limited factor, especially when the fuel cell operates at low current densities. Thus, the reduction of  $\text{H}_2$  crossover when a thinner membrane is used is a challenge which necessitates the investigation of hydrogen crossover, especially when the fuel cell is operated at higher temperatures at which the  $\text{H}_2$  crossover is more severe.

Hydrogen crossover is an undesirable diffusion of hydrogen from the anode to the cathode through the membrane. Hydrogen crossover can have at least three effects, including fuel efficiency reduction, cathode potential depression, and aggressive peroxide radical formation. The hydrogen which crosses over can directly react with oxygen at the cathode surface, resulting in a lower cathode potential than that of a lower fuel cell (OCV), a result confirmed by our recent study of temperature-dependent OCV in PEM fuel cells [5]. More severely, this direct reaction between  $\text{H}_2$  and  $\text{O}_2$  at the cathode can produce peroxide radicals, which not only attack the catalyst layer but also the membrane, causing significant catalyst-layer and membrane degradation [6]. In addition, it has been confirmed that the formation of hot-points [7] or hydrogen peroxide [8,9] by the highly exothermal chemical reaction between  $\text{H}_2$  and  $\text{O}_2$  can also lead to pin-holes in membranes, destroying the MEA and causing safety problems. An accelerated sintering of catalysts could be also caused by this hydrogen crossover [6]. Therefore, the measurement of hydrogen crossover, in particular at OCV, at which the most severe crossover occurs, is of importance for the fundamental understanding and practical mitigation of fuel cell degradation and membrane failure.

Since  $\text{H}_2$  crossover is a diffusion-controlled process, the PEM structure [10] and fuel cell operating conditions [8,11] can greatly influence the crossover process. For example, it can be expected that the uses of thinner membranes and the operation of fuel cells at high temperature and high pressure will facilitate  $\text{H}_2$  crossover.  $\text{H}_2$  crossover at temperatures up to 80  $^\circ\text{C}$  has been investigated by a method of limiting current, primarily using linear sweep voltammetry (LSV) through various membranes: Nafion [12–15], Nafion-palladium [12], sulfonated poly(ether ether ketone) (SPEEK) [16], and Flemion [17]. It was found that the water content in the membrane [12] and gas pressure [8] had a great effect on the  $\text{H}_2$  crossover rate. An in situ method for limiting  $\text{H}_2$  crossover current measurements in subscale and full-size single cells has also been developed [18].

High-temperature PEM fuel cells (>80  $^\circ\text{C}$ ) have several advantages over those operated at lower temperatures [19]. However, the data available for  $\text{H}_2$  crossover at temperatures higher than 80  $^\circ\text{C}$  are very limited. It can be expected that the problems associated with high-temperature  $\text{H}_2$  crossover would be worse than those at lower temperatures. In a recent model study based on previously developed models [20,21], Rama et al. [22] discussed the effects of membrane thickness, operating pressure and temperature, and current density on  $\text{H}_2$  crossover using a concentrated solution membrane system containing four species: water, electrolyte membrane, proton and hydrogen. The simulated results indicate that the increase in the  $\text{H}_2$  crossover found when temperatures rose from 80 to 100  $^\circ\text{C}$  at 3 atm was

attributed to the dependence of the  $\text{H}_2$  diffusion coefficient on temperature [22]. The use of Nafion-based composite membranes could reduce  $\text{H}_2$  crossover with a more crystalline region being provided after heat-treating at 120  $^\circ\text{C}$  [23] or with less porosity resulting from well-dispersed small additive particles [24].

The structure and chemistry of PEM are critical in the control of proton and water transport and  $\text{H}_2$  crossover. The three-region model [25] and cluster-network model [26] have been widely used to describe the microstructure of perfluorosulfonated membranes such as Nafion-type membranes. However, there is still no agreement as to whether the gas permeates through Nafion in the intermediate region, which consists mainly of the flexible amorphous part of the perfluorocarbon backbone [13,14], or through the hydrated ionic clusters, which contain mainly water and sulfonic acid groups [12]. A compromise solution, that the gas permeates through both regions, has also been suggested [27].

In order to improve our understanding of the  $\text{H}_2$  crossover phenomena, it is necessary to obtain more information about  $\text{H}_2$  crossover at elevated temperatures and to further clarify the nature of  $\text{H}_2$  crossover from the anode to the cathode via the MEA. This investigation was carried out to examine the effects of temperature, pressure, and relative humidity on  $\text{H}_2$  crossover. A multilayer diffusion model was proposed to describe  $\text{H}_2$  crossover through an MEA. Some  $\text{H}_2$  crossover parameters were evaluated with a focus on the high-temperature  $\text{H}_2$  crossover behaviour observed in a fuel cell assembling environment.

## 2. Experimental

The MEA, with an active area of 4.4  $\text{cm}^2$ , was prepared by hot pressing the anode, a Nafion 112 (or 117) membrane, and cathode together at 135  $^\circ\text{C}$  and 75  $\text{kg cm}^{-2}$  for 2 min. The gas diffusion electrode (GDE) was prepared by spraying a homogeneous catalyst ink, consisting of catalyst, Nafion solution, and *iso*-propanol, onto a gas diffusion layer (GDL). This GDL was a PTFE- and carbon-black impregnated carbon paper (Toray, TGP-H-060). E-Tek 20% Pt/Ru/C and 40% Pt/C were used as the anode and cathode catalysts, with a total Pt loading of  $\sim 1.0 \text{ mg cm}^{-2}$ . The total Nafion loading in the MEA was 1.4  $\text{mg cm}^{-2}$ . The Nafion 112 and 117 (DuPont) used for the PEM were treated in 3%  $\text{H}_2\text{O}_2$  (aq), 1 M  $\text{H}_2\text{SO}_4$  (aq); rinsed; and then soaked separately in deionized water for 1 h at 60–80  $^\circ\text{C}$ , followed by a careful washing with double-distilled water.

An in-house single fuel cell hardware, described in our previous paper [3], was used for measuring fuel cell performance and  $\text{H}_2$  crossover. A bladder pressure of 4.4 atm was used to hold the single cell together and provide sufficient electrical contact between the MEA and the graphite bipolar plates. Both graphite plates had the same serpentine flow fields. A 100 W Fideris fuel cell test station controlled by FC Power software and equipped with an in-house modified humidifier which could reach 100% RH at 120  $^\circ\text{C}$  was used for fuel cell polarization at different temperatures, pressures, and relative humidities. When the cell was controlled at a certain RH, the  $\text{H}_2$  and air (or nitrogen), before they were fed into the anode and cathode, were first

passed through their corresponding humidifiers, in which they were humidified at the same temperature as that of the fuel cell. During all tests, the pressures for both the anode and the cathode sides were kept at the same level.

Before H<sub>2</sub> crossover was measured at various conditions, the fuel cell assembled with an MEA was conditioned at 80 °C, 3 atm, and 100% RH with a load of 1.0 A cm<sup>-2</sup> for 4 h, followed by a collection of steady-state polarization data in the current density range of 0–2.5 A cm<sup>-2</sup>. After that, a humidified nitrogen stream was introduced into the cathode to remove the air. After 30 min of N<sub>2</sub> flushing through the cathode, a Solartron 1287 potentiostat was connected to the fuel cell for H<sub>2</sub> crossover measurements, with the working electrode probe connected to the cathode and the counter/reference electrode probes connected together to the anode. A steady-state electrochemical method was used to record the current produced from the oxidation of crossed H<sub>2</sub> from the anode at an applied cathode potential of 0.5 V relative to the potential of the H<sub>2</sub>-flushed anode. At this cathode potential, all H<sub>2</sub> that has crossed over from the anode to the cathode should be completely oxidized, giving a current indicative of the amount of hydrogen that has crossed over. The obtained current was defined as  $I_{H_2}^{cross}$  (unit: A) for the calculation of H<sub>2</sub> crossover rate ( $J_{H_2}^{cross}$ , mol cm<sup>-2</sup> s<sup>-1</sup>). The following Faraday's equation was used to obtain the H<sub>2</sub> crossover rate:

$$J_{H_2}^{cross} = \frac{I_{H_2}^{cross}}{nFA} \quad (2)$$

where  $n$  is the electron number of H<sub>2</sub> oxidation (=2),  $F$  the Faraday constant (A s mol<sup>-1</sup>), and  $A$  is the MEA active area (cm<sup>2</sup>). It is worthwhile to point out that the H<sub>2</sub> crossover rate measured by this procedure should represent the case at OCV conditions rather than those at load.

In order to obtain the dependencies of the H<sub>2</sub> permeability coefficient ( $\psi_{H_2}^{PEM}$ ) on the temperature ( $T$ ), fuel cell backpressure ( $P$ ), and relative humidity (RH), the H<sub>2</sub> partial pressure ( $P_{H_2}^a$ ) as a function of  $T$ ,  $P$ , and RH must be known. The values of the H<sub>2</sub> partial pressure were calculated based on the fuel cell anode inlet total pressure ( $P_{inlet}^a$ ), the H<sub>2</sub> partial pressure in the inlet stream ( $P_{H_2-inlet}^a$ ), and the outlet total pressure, which is equal to the backpressure ( $P$ ). The following Eq. (3) was used to calculate the partial pressure inside the anode chamber by assuming that the pressure drop along the superfine channel is linear:

$$P_{H_2}^a = \left( \frac{P_{H_2-inlet}^a}{P_{inlet}^a} \right) \left( \frac{P_{inlet}^a + P}{2} \right) \quad (3)$$

where  $P_{inlet}^a$  and  $P$  can be measured experimentally, and  $P_{H_2-inlet}^a$  can be obtained based on the inlet H<sub>2</sub> stream flow rate, the water vapour partial pressure, and the inlet total pressure ( $P_{inlet}^a$ ). Note that the amount of H<sub>2</sub> that crosses over from the anode to the cathode during the measurements is expected to be too small to cause a significant drop in the H<sub>2</sub> partial pressure. Therefore, it is not necessary to make a crossover correction to the value of  $P_{H_2}^a$ . The H<sub>2</sub> partial pressure data obtained according to Eq. (3) at various  $T$ s,  $P$ s, and RHs were used to calculate H<sub>2</sub> permeability coefficients, as discussed in later sections of this paper.

### 3. Results and discussion

#### 3.1. Theoretical description of the H<sub>2</sub> crossover through the MEA

In a fuel cell, hydrogen crossover usually takes place from the anode side to the cathode side through a MEA, as illustrated in Fig. 1, which schematically represents a typical MEA structure consisting of five layers: anode gas diffusion layer (GDL(a)), anode catalyst layer (CL(a)), proton exchange membrane (PEM), cathode CL (CL(c)), and cathode GDL (GDL(c)). By Fick's first law, the steady-state H<sub>2</sub> crossover rate can be treated as the diffusion through the MEA:

$$J_{H_2}^{cross} = \left( \frac{D_{H_2}}{l} \right) (C_{H_2}^a - C_{H_2}^c) \quad (4)$$

where  $C_{H_2}^a$  and  $C_{H_2}^c$  are the concentration of H<sub>2</sub> at the interface between the anode H<sub>2</sub> gas phase and the GDL(a) surface, and the concentration at the interface between the cathode air (or nitrogen) gas phase and the GDL(c) surface, respectively.  $D_H$  is the overall diffusion coefficient and  $l$  is the MEA thickness.

As depicted in Fig. 1, the overall resistance for H<sub>2</sub> diffusion ( $l/D_{H_2}$ ) through an MEA is the sum of the resistance for H<sub>2</sub> diffusion through each layer. Therefore,  $D_{H_2}$  can be expressed as Eq. (5) based on the equivalent resistance of the five layers in series:

$$\frac{l}{D_{H_2}} = \frac{l_{GDL}^a}{D_{GDL}^a} + \frac{l_{CL}^a}{D_{CL}^a} + \frac{l_{PEM}}{D_{PEM}} + \frac{l_{CL}^c}{D_{CL}^c} + \frac{l_{GDL}^c}{D_{GDL}^c} \quad (5)$$

where  $l_{GDL}^a$ ,  $l_{CL}^a$ ,  $l_{PEM}$ ,  $l_{CL}^c$ , and  $l_{GDL}^c$  are the thicknesses of GDL(a), CL(a), PEM, CL(c), and GDL(c), respectively, and  $l$  can be expressed as Eq. (6):

$$l = l_{GDL}^a + l_{CL}^a + l_{PEM} + l_{CL}^c + l_{GDL}^c \quad (6)$$

In (5)  $D_{GDL}^a$ ,  $D_{CL}^a$ ,  $D_{PEM}$ ,  $D_{CL}^c$ , and  $D_{GDL}^c$  are the diffusion coefficients of H<sub>2</sub> in GDL(a), CL(a), PEM, CL(c), and GDL(c), respectively. Hence,  $D_{H_2}$  can be deduced from the rearrange-

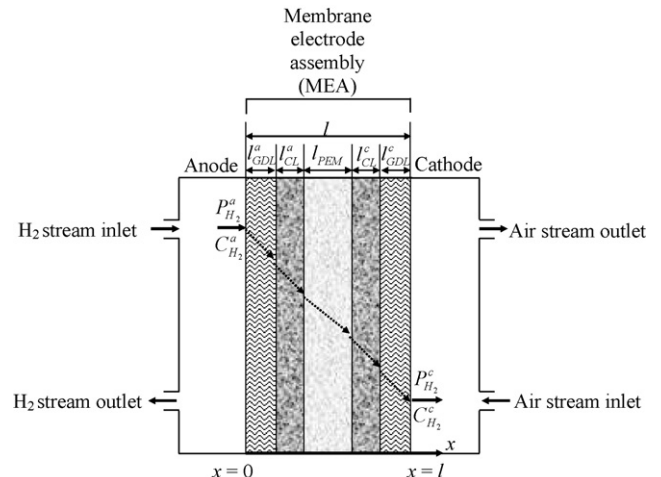


Fig. 1. Schematic diagram of hydrogen crossover through the (MEA) in a fuel cell.

ment of Eq. (5):

$$D_{H_2} = \frac{l_{GDL}^a + l_{CL}^a + l_{PEM} + l_{CL}^c + l_{GDL}^c}{(l_{GDL}^a/D_{GDL}^a) + (l_{CL}^a/D_{CL}^a) + (l_{PEM}/D_{PEM}) + (l_{CL}^c/D_{CL}^c) + (l_{GDL}^c/D_{GDL}^c)} \quad (7)$$

Eq. (7) clearly indicates that the diffusion from every layer depicted in Fig. 1 contributes to the overall  $H_2$  diffusion. Since  $H_2$  diffused from the anode to the cathode would fully react with  $O_2$  at the interface of PEM/CL(c), or would be completely oxidized electrochemically by the measurement procedure described in Section 2, the  $H_2$  diffusion paths  $l_{CL}^c$  and  $l_{GDL}^c$  are equal to zero. Thus, Eq. (7) can be simplified as Eq. (8):

$$D_{H_2} = \frac{l_{GDL}^a + l_{CL}^a + l_{PEM}}{(l_{GDL}^a/D_{GDL}^a) + (l_{CL}^a/D_{CL}^a) + (l_{PEM}/D_{PEM})} \quad (8)$$

For the anode GDL and CL, the diffusion coefficients  $D_{GDL}^a$  and  $D_{CL}^a$  can be expressed as Eqs. (9) and (10), respectively, based on the Bruggeman correlation [28,29]

$$D_{GDL}^a = D_{H_2}^g [\varepsilon_{GDL}^a (1 - s_{GDL}^a)]^\tau \quad (9)$$

$$D_{CL}^a = D_{H_2}^g [\varepsilon_{CL}^a (1 - s_{CL}^a)]^\tau \quad (10)$$

where  $D_{H_2}^g$  is the vapour-phase  $H_2$  diffusion coefficient; its typical value at  $80^\circ C$  is  $2.63 \times 10^{-2} \text{ cm}^2 \text{ s}^{-1}$  [28,29].  $\varepsilon_{GDL}^a$  and  $\varepsilon_{CL}^a$  are the porosities of GDL and CL, respectively, and  $s_{GDL}^a$  and  $s_{CL}^a$  are the water saturation degrees inside GDL and CL, respectively. Their typical values are in the range of 0–1.  $\tau$  is the tortuosity, which is often assumed to be 1.5.

Although  $D_{PEM}$  for Nafion membranes has been formulated empirically as a function of temperature ( $T$ ) in the literature [30], the backpressure ( $P$ ) and RH dependences measured in a fuel cell environment should be closer to a real situation.

Based on the literature, the typical thicknesses of PEMs are  $\sim 50 \mu\text{m}$  for Nafion 112 and  $\sim 175 \mu\text{m}$  for Nafion 117. The thickness for GDL is about  $200 \mu\text{m}$ . The diffusion coefficients of  $H_2$  crossover through GDL(a), CL(a), and PEM layers for a conventional MEA in a fuel cell at  $80^\circ C$  are at the orders of  $\sim 10^{-2} \text{ cm}^2 \text{ s}^{-1}$ ,  $\sim 10^{-4} \text{ cm}^2 \text{ s}^{-1}$ , and  $\sim 10^{-6} \text{ cm}^2 \text{ s}^{-1}$ , respectively. Combining these values into Eq. (8), it can be seen that the  $H_2$  crossover through an MEA is predominantly caused by the membrane diffusion, which is mainly due to the smallest  $D_{PEM}$  value. The contributions from both GDL and CL diffusions can be ignored. Therefore, Eq. (8) can be simplified further to be Eq. (11) by eliminating the contributions of GDL(a), and CL(a):

$$\frac{l}{D_{H_2}} = \frac{l_{PEM}}{D_{PEM}} \quad (11)$$

Eq. (11) indicates that the overall diffusion coefficient value of  $H_2$  crossover (or diffusion) through the MEA, obtained experimentally, is equal to that through the PEM. This has been recognized widely in the literature. For example, the dependence of  $H_2$  crossover on the membrane thickness was confirmed by Kocho et al. [15], who compared the  $H_2$  crossover currents measured in Nafion 112 and in Nafion 111 in the temperature range

of  $25\text{--}80^\circ C$ , suggesting that the dominant factor limiting the  $H_2$  crossover rate is the membrane diffusion coefficient.

Combining Eq. (11) with Eq. (4), the  $H_2$  crossover rate can be expressed as Eq. (12) if the crossed  $H_2$  at the cathode can be completely oxidized ( $C_{H_2}^c = 0$ ):

$$J_{H_2}^{\text{cross}} = \left( \frac{D_{PEM}}{l_{PEM}} \right) C_{H_2}^a \quad (12)$$

Eq. (12) suggests that the  $H_2$  crossover rate through the MEA is equal to that through the PEM. An alternative expression for Eq. (12) can be introduced as Eq. (13):

$$J_{H_2}^{\text{cross}} = \left( \frac{K_{H_2}^{\text{PEM}} D_{PEM}}{l_{PEM}} \right) P_{H_2}^a \quad (13)$$

where  $K_{H_2}^{\text{PEM}}$  is a  $H_2$  partial pressure-related solubility coefficient in the PEM with a unit of  $\text{mol cm}^{-3} \text{ atm}^{-1}$ , and  $P_{H_2}^a$  is the  $H_2$  partial pressure in the anode feed stream with a unit of atm. The product of the solubility ( $K_{H_2}^{\text{PEM}}$ ) and diffusivity ( $D_{PEM}$ ) can be defined as the permeability coefficient of  $H_2$  in the PEM,  $\psi_{H_2}^{\text{PEM}}$ , with a unit of  $\text{mol cm}^{-1} \text{ atm}^{-1} \text{ s}^{-1}$ :

$$\psi_{H_2}^{\text{PEM}} = K_{H_2}^{\text{PEM}} D_{PEM} \quad (14)$$

Eq. (15) is obtained by substituting Eq. (14) into Eq. (13):

$$J_{H_2}^{\text{cross}} = \left( \frac{\psi_{H_2}^{\text{PEM}}}{l_{PEM}} \right) P_{H_2}^a \quad (15)$$

Similar to  $D_{H_2}$  (or  $D_{PEM}$ ),  $\psi_{H_2}^{\text{PEM}}$  is also a function of  $T$ ,  $P$ , RH, water content, and the nature of the membrane [21]. For a defined fuel cell system, the value of  $l_{PEM}$  in Eq. (15) does not vary substantially; the other two factors,  $\psi_{H_2}^{\text{PEM}}$  and  $P_{H_2}^a$ , may become more dominate to influence  $H_2$  crossover, as suggested by Eq. (15).

The permeability coefficient,  $\psi_{H_2}$ , can then be determined by rearranging Eq. (15):

$$\psi_{H_2}^{\text{PEM}} = \frac{J_{H_2}^{\text{cross}} l_{PEM}}{P_{H_2}^a} \quad (16)$$

In Eq. (16), the values of  $J_{H_2}^{\text{cross}}$  and  $P_{H_2}^a$  can be measured experimentally, and then  $\psi_{H_2}^{\text{PEM}}$  can be calculated at different  $T$ s,  $P$ s, and RHs if  $l_{PEM}$  is known (for Nafion 112,  $l_{PEM} = 50 \mu\text{m}$ ; for Nafion 117,  $l_{PEM} = 175 \mu\text{m}$ ).

### 3.2. $H_2$ crossover rate as a function of $T$ , $P$ , and RH

As described in Section 2, hydrogen crossover was measured using a steady-state electrochemical method. The current induced by crossed-over hydrogen ( $J_{H_2}^{\text{cross}}$ ) was used to simulate the values of  $\psi_{H_2}^{\text{PEM}}$  at various conditions. Table 1 lists the  $H_2$  crossover rates obtained at various  $T$ s,  $P$ s, and RHs using a Nafion 112-based MEA. In general, the magnitude of  $J_{H_2}^{\text{cross}}$  goes up monotonically with increases in temperature and backpressure. However, the RH effect is more complicated than those of temperature and backpressure.



Table 1

Measured H<sub>2</sub> crossover rate at various temperatures, backpressures, and relative humidities using a Nafion 112-based MEA

Temperature (°C)	RH (%)	Measured H <sub>2</sub> crossover rate (mol cm <sup>-2</sup> s <sup>-1</sup> )		
		3.04 <sup>a</sup>	2.02 <sup>a</sup>	1.00 <sup>a</sup>
80	100	2.04E-08	1.30E-08	3.78E-09
	70	1.91E-08	1.20E-08	4.14E-09
	50	1.80E-08	1.16E-08	4.51E-09
	25	1.74E-08	1.14E-08	5.30E-09
100	100	2.69E-08	1.48E-08	4.13E-09
	70	2.82E-08	1.43E-08	5.91E-09
	50	2.57E-08	1.35E-08	7.00E-09
	25	2.23E-08	1.42E-08	8.23E-09
120	100	3.05E-08	1.65E-08	8.13E-09
	70	4.16E-08	2.33E-08	1.19E-08
	50	5.02E-08	2.90E-08	1.61E-08
	25	6.28E-08	4.26E-08	2.17E-08

The thickness of the Nafion 112 was adopted as 50 μm.

<sup>a</sup> Backpressure (atm).

### 3.2.1. Temperature effect on the H<sub>2</sub> crossover rate

All the data listed in Table 1 follow a trend in which the H<sub>2</sub> crossover rate increases with increasing temperature. This observation has also been reported in the temperature range of 60–80 °C, and was explained according to the increase in membrane flexibility when the temperature and humidity were increased [8]. Because the permeability coefficient,  $\psi_{\text{H}_2}^{\text{PEM}}$ , is directly related to the H<sub>2</sub> solubility and diffusion coefficients [15], any effect of temperature on these two coefficients could cause a change of  $\psi_{\text{H}_2}^{\text{PEM}}$ . The increase of  $\psi_{\text{H}_2}^{\text{PEM}}$  with increasing temperature probably indicates that the effect of temperature on the diffusion coefficient is more significant than that on the solubility coefficient.

### 3.2.2. Backpressure effect on the H<sub>2</sub> crossover rate

Fig. 2 shows the H<sub>2</sub> crossover rate as a function of backpressure at 50% RH and three different temperatures. A monotonic increase in the H<sub>2</sub> crossover rate with increasing backpressure can be observed. It is understandable that an increase in back-

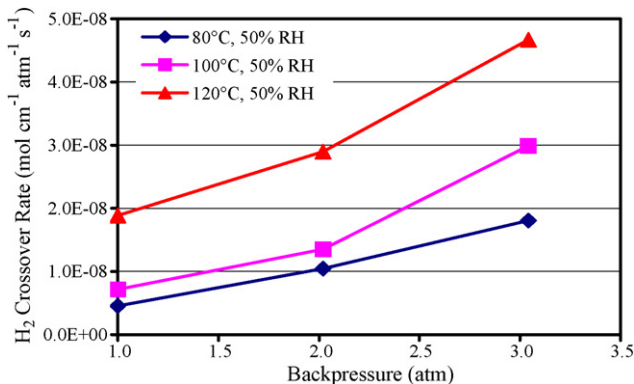


Fig. 2. H<sub>2</sub> crossover rate as a function of fuel cell backpressure at 50% relative humidity and three different temperatures as marked in the figure. Nafion 112-based MEA with an active area of 4.4 cm<sup>2</sup>. Anode H<sub>2</sub> stream flowrate: 0.1 L min<sup>-1</sup>; cathode N<sub>2</sub> stream flowrate: 0.5 L min<sup>-1</sup>. Cathode potential: 0.5 V vs. anode hydrogen electrode.

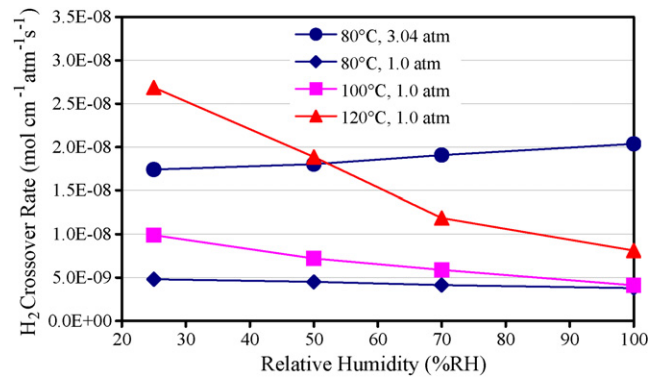


Fig. 3. H<sub>2</sub> crossover rate as a function of relative humidity in fuel cell anode feed stream at three different temperatures as marked in the figure. Nafion 112-based MEA with an active area of 4.4 cm<sup>2</sup>. Anode H<sub>2</sub> stream flowrate: 0.1 L min<sup>-1</sup>; cathode N<sub>2</sub> stream flowrate: 0.5 L min<sup>-1</sup>. Cathode potential: 0.5 V vs. anode hydrogen electrode.

pressure will result in a H<sub>2</sub> partial pressure increase, which then creates a larger pressure difference across the PEM. As expected by Eq. (15), the H<sub>2</sub> crossover rate,  $J_{\text{H}_2}^{\text{cross}}$ , will be increased if the H<sub>2</sub> partial pressure,  $P_{\text{H}_2}^a$ , increases.

### 3.2.3. RH effect on the H<sub>2</sub> crossover rate

Fig. 3 shows the typical data obtained regarding the RH effect on the H<sub>2</sub> crossover rate. It seems that the RH effect is more complicated than those of temperature and backpressure. For example, at 80 °C and 3.0 atm, the H<sub>2</sub> crossover rate increases as the RH increases, while at the same temperature and a backpressure of 1.0 atm, the RH increase reduces the crossover rate. However, at high temperatures (120 °C), the effect of RH on H<sub>2</sub> crossover is always negative; that is, the H<sub>2</sub> crossover is reduced with increasing RH.

For a fuel cell anode feed stream, the inlet total pressure is the sum of the H<sub>2</sub> partial pressure and the water vapour partial pressure ( $P_{\text{H}_2\text{O}}^a$ ):

$$P_{\text{H}_2\text{-inlet}}^a = P_{\text{H}_2}^a + P_{\text{H}_2\text{O}}^a \quad (17)$$

For a controlled total pressure, an increase in the RH level results in an increase in the magnitude of  $P_{\text{H}_2\text{O}}^a$ . Consequently, the  $P_{\text{H}_2}^a$  is reduced accordingly, as expected from Eq. (17). It can be seen from Eq. (15) that the H<sub>2</sub> crossover will be reduced if  $P_{\text{H}_2}^a$  is decreased, which is the case at 100 °C and 120 °C, as shown in Fig. 3. However, the increasing RH level will increase the water content in the membrane, which may cause changes in both the value of H<sub>2</sub> solubility and the diffusion coefficients. At lower temperature (80 °C), the value increase in the permeability coefficient caused by an RH increase could be greater than that caused by a  $P_{\text{H}_2}^a$  reduction, resulting in the increasing trend observed in Fig. 3 (80 °C, 3.04 atm).

## 3.3. H<sub>2</sub> permeability coefficient as a function of T, P, and RH

The data in Table 2 were obtained using Eq. (16) based on the measured data of H<sub>2</sub> crossover rate and partial pressure. Table 2 indicates that an increase in temperature can effectively increase

Table 2

Calculated H<sub>2</sub> permeability coefficients at various temperatures, backpressures, and relative humidities used a Nafion 112-based MEA

Temperature (°C)	RH (%)	H <sub>2</sub> permeability coefficient (mol cm <sup>-1</sup> atm <sup>-1</sup> s <sup>-1</sup> )		
		3.04 <sup>a</sup>	2.02 <sup>a</sup>	1.00 <sup>a</sup>
80	100	3.87E-11	3.71E-11	1.88E-11
	70	3.47E-11	3.24E-11	1.93E-11
	50	3.10E-11	2.95E-11	2.06E-11
	25	2.93E-11	2.80E-11	2.21E-11
100	100	6.13E-11	5.24E-11	2.40E-11
	70	5.74E-11	4.42E-11	3.11E-11
	50	4.89E-11	3.84E-11	3.39E-11
	25	3.90E-11	3.64E-11	3.54E-11
120	100	1.04E-10	9.14E-11	6.94E-11
	70	1.08E-10	9.52E-11	7.94E-11
	50	1.12E-10	9.80E-11	8.95E-11
	25	1.18E-10	1.08E-10	1.01E-10

The thickness of the Nafion 112 was adopted as 50 μm.

<sup>a</sup> Backpressure (atm).

the H<sub>2</sub> permeability coefficient at any given backpressure and RH. This trend was also observed by Kocha et al. [15]. Table 2 also shows that the temperature effect is much larger than those affected by backpressure and RH.

The temperature dependence of the permeability coefficient can be used to obtain the activation energy of H<sub>2</sub> permeability ( $E_{H_2}^{PEM}$ ) through an Arrhenius relation:

$$\psi_{H_2}^{PEM} = \psi_{H_2}^0 e^{-E_{H_2}^{PEM}/RT} \quad (18)$$

where  $\psi_{H_2}^0$  is the maximum H<sub>2</sub> permeability coefficient (i.e., at infinite temperature), and  $R$  is the gas constant. Eq. (18) can be written into a logarithmic form:

$$\ln \psi_{H_2}^{PEM} = \ln \psi_{H_2}^0 + \left( -\frac{E_{H_2}^{PEM}}{R} \right) \frac{1}{T} \quad (19)$$

Thus, a straight line is expected by plotting  $\ln \psi_{H_2}^{PEM}$  against  $1/T$ , and the values of  $\psi_{H_2}^0$  and  $E_{H_2}^{PEM}$  can be evaluated from the intercept and slope of the plot, respectively. Fig. 4 gives the typical Arrhenius plots for both Nafion 112 and Nafion 117-based MEAs at 3.04 atm backpressure and 100% RH in the temperature range of 40–120 °C. Both Nafion 112 and Nafion 117 give similar values of activation energy and the maximum permeability coefficient, which are 24.7 kJ mol<sup>-1</sup> and  $1.39 \times 10^{-7}$  mol cm<sup>-1</sup> atm<sup>-1</sup> s<sup>-1</sup>, respectively. The value of  $E_{H_2}^{PEM}$  is close to the value of 21.03 kJ mol<sup>-1</sup> reported in the literature using Nafion 111 and Nafion 112 membranes [15]. The fact that the values of activation energy and maximum permeability obtained from Nafion 112 are similar to those obtained from Nafion 117 strongly suggests that the membrane thickness is not the dominating factor affecting the H<sub>2</sub> permeability coefficient. Yoshida et al. [20] also found that the  $\psi_{H_2}^{PEM}$  values for three types of Flemion membranes and Nafion 117 were nearly the same and were independent of the membrane thickness in the region from 15 to 230 μm at 70 °C. However, at 120 °C, the  $\psi_{H_2}^{PEM}$  values were an order larger than those at 30 °C. Hence,

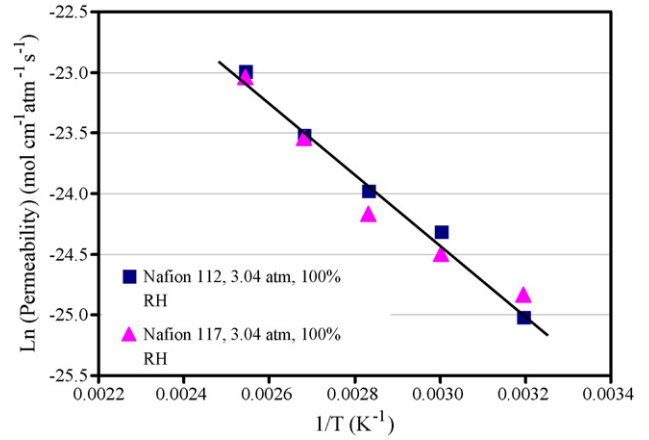


Fig. 4. Arrhenius plots of H<sub>2</sub> permeability coefficient for both Nafion 112 and Nafion 117-based MEAs at 3.04 atm backpressure and 100% RH in the temperature range of 40–120 °C. Both MEAs have an active area of 4.4 cm<sup>2</sup>. Anode H<sub>2</sub> stream flowrate: 0.1 L min<sup>-1</sup>; cathode N<sub>2</sub> stream flowrate: 0.5 L min<sup>-1</sup>. Cathode potential: 0.5 V vs. anode hydrogen electrode.

the H<sub>2</sub> permeation coefficient is more strongly dependent on temperature than on membrane thickness.

Therefore, Eq. (18) can be rewritten as Eq. (20) for the operation conditions of backpressure 3.04 atm and 100% RH:

$$\psi_{H_2}^{PEM} = 1.39 \times 10^{-7} e^{-2980/T} \quad (20)$$

Table 2 shows that the permeability coefficient is also dependent on the backpressure and RH. Similar to the effect of temperature, increasing backpressure can always raise the value of the H<sub>2</sub> permeability coefficient at any given RH. However, the effect of RH on the permeability coefficient seems to be more complicated. For example, the trends obtained when RH is increased at lower temperatures (80 °C and 100 °C) and higher backpressures (2.02 atm and 3.04 atm) are opposite to those obtained at higher a temperature (120 °C) and lower pressure (1.0 atm). The change in the water content of the membrane at different backpressures and temperatures can affect both H<sub>2</sub> solubility and diffusivity, resulting in the change in the permeability coefficient. More work has been scheduled to investigate the mechanism of H<sub>2</sub> crossover through different membranes in order to improve our fundamental understanding of these effects.

For a simple approach, the effects of backpressure and RH were empirically simulated and can be expressed as Eq. (21):

$$\psi_{H_2}^{PEM} = 1.39 \times 10^{-7} e^{-(2980f(P,RH))/T} \quad (21)$$

where  $f(P, RH)$  reflects the effects of backpressure and RH on the activation energy. However, in order to obtain the exact expression for this  $f(P, RH)$  by empirical simulation, a statistically significant amount of data is necessary, which is beyond the scope of this paper.

#### 4. Conclusions

The hydrogen crossover in a PEM fuel cell environment was measured experimentally in the temperature range of 40–120 °C, focusing especially on the higher temperature range of 80–120 °C. An H<sub>2</sub> permeation model based on an MEA con-

sisting of five layers – anode GDL, anode CL, PEM, cathode GDL, and cathode CL – was constructed to obtain a simple expression for the calculation of H<sub>2</sub> permeability coefficients based on the measured H<sub>2</sub> crossover rates and controlled H<sub>2</sub> partial pressures at various temperatures, backpressures, and relative humidities. The model analysis suggests that the H<sub>2</sub> permeation through the PEM is a limiting step in the overall process.

The H<sub>2</sub> permeability coefficients obtained as a function of temperature, backpressure, and RH show that an increase in temperature can effectively increase the H<sub>2</sub> permeability coefficient at any given backpressure and RH. Similarly, increasing backpressure can always increase the value of the H<sub>2</sub> permeability coefficient at any given RH. However, the effect of RH on the permeability coefficient seems to be more complicated. The temperature effect is much larger than those caused by backpressure and RH. An equation to describe the *T* dependencies was obtained based on experimental data by simulation.

The activation energies and maximum permeability coefficients of H<sub>2</sub> permeation through both Nafion 112 and Nafion 117-based MEAs were evaluated separately. Both MEAs showed similar values for permeability coefficients and activation energies, suggesting that the membrane thickness has no significant effect on the H<sub>2</sub> permeation mechanism.

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